Design of organic molecules and molecular assemblies with electron spin alignment

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Abstract - The electron spin-spin interaction between the two triplet diphenylcarbene units incorporated in the [2.2]paracyclophane skeleton has been investigated by ESR spectroscopy. The pseudoortho and pseudopara bis(phenylmethylenyl)[2.2]paracyclophanes are in the ground quintet state, while the pseudometa isomer is singlet. A conformational analysis of these high-spin dicarbenes has been carried out by examining their dipolar interaction tensors. The McConnell's theory on the ferromagnetic interaction between two organic radicals has been demonstrated and a new strategy for increasing the dimension of high-spin aromatic molecules was found. On the basis of these findings, the molecular aggregates of some p-alkoxy substituted diphenylcarbenes were generated in the crystals of the precursor diazodiphenylmethanes to give 4 to 8 spins aligned among the neighboring triplet molecules.

INTRODUCTION

The large majority of organic molecules have closed shell electronic structure, i.e., they have singlet ground state with equal number of electrons having α and β spin. Most organic compounds are therefore diamagnetic. There are some organic molecules that have open-shell structures where not all electrons are paired. Many of these molecules have one odd electron and therefore doublet ground state. These free radicals are paramagnetic. The alignment of spins in molecular systems becomes an issue when there is an interaction between two doublet centers as in diradicals and radical pairs. The Coulombic repulsion between electrons lifts the zeroth-order degeneracy and gives rise to singlet and triplet states of different total energy for these chemical entities. It is highly important in the chemistry of diradicals (ref. 1) and radical pairs (ref. 2) which state lies lower in energy. The conditions for having the ground triplet states and the importance of the singlet-triplet radical pairs in free radical reactions have been amply discussed.

While diradicals and radical pairs are accepted as important reaction intermediates, the chemistry of these and more highly spin-aligned systems is considered also to be the best area where one can look for interesting magnetic properties from organic materials. Recently we have shown that polycarbenes 1 and 2 generated by photolysis of the corresponding polydiazo compounds have their highest possible spin multiplicities in the ground states (ref. 3). The magnetic susceptibility measurement has revealed that nonet tetracarbene 1(n = 4) has a paramagnetic property of very large magnetic moments (S = 8/2)(ref. 4). Higher homologs of 1 and 2 are expected to be best characterized as organic molecular superparamagnets (ref. 5). Curie plots of the paramagnetic susceptibility of 1(n = 4) in 2-methyltetrahydrofuran glass or in polycrystalline sample have shown, however, that the intermolecular interaction in the fortuitously formed aggregates under these conditions is antiferromagnetic. It appears to be the rule rather than exception that the intermolecular interaction between open-shell organic

molecules should be antiferromagnetic. Therefore, it is imperative for obtaining the spin alignment in macroscopic scale throughout the molecular assembly to design systems in which molecular interaction could become ferromagnetic. There appear to be a couple of plausible methods for achieving this objective. One is to take advantage of a weak hyperconjugative interaction through a short chain connecting high-spin molecules (ref. 6). The other is a design of molecular stacking of spin-distributed aromatic rings leading to intermolecular spin alignment (ref. 7). In this paper, we would like to discuss the latter strategy on the basis of the McConnell's theory (ref. 8) and show a partial success of the application of this method.

DESIGN OF A MODEL FOR FERROMAGNETIC MOLECULAR STACKING OF DIPHENYLCARBENE UNITS

Aromatic molecules have a tendency to stack in organized systems. It would therefore be desirable if requirements for parallel spins, i.e., ferromagnetic stacking between two aromatic radicals are found. For a set of carbenes A and B each initially in a triplet state, the phenomenological interaction Hamiltonian is written out as

$$\hat{\mathbf{H}}^{AB} = -2 \sum_{\mathbf{m}, \mathbf{n} = 1, 2} J_{\mathbf{m}\mathbf{n}}^{AB} \mathbf{S}_{\mathbf{m}}^{A} \cdot \mathbf{S}_{\mathbf{n}}^{B} \tag{1}$$

where J_{mn}^{AB} is the valence-bond exchange integral between the odd electrons m and n on carbenes A and B, respectively, with respective spin angular momentum operators \mathbf{S}_{m}^{A} and \mathbf{S}_{n}^{B} . For weak π -type interactions, we express the operator $\mathbf{S}^{A} \cdot \mathbf{S}^{B}$ as

$$-2J^{AB}\mathbf{S}^{A} \cdot \dot{\mathbf{S}}^{B} = -J^{AB}(\mathbf{S}^{2} - \mathbf{S}^{A2} - \mathbf{S}^{B2})$$
(2)

where $\mathbf{S} = \mathbf{S}^A + \mathbf{S}^B$ is the operator for the total spin angular momentum, with eigenvalue S. For two spins \mathbf{S}^A and \mathbf{S}^B with eigenvalues $S^A = 1$ and $S^B = 1$, the expectation value of $\hat{\mathbf{H}}$ is

$$E = -J[S(S+1) - 4]$$
 (3)

The state of eigenvalue -2J has five-fold degeneracy with a total spin S=2, which is a quintet state. The eigenfunctions of the quintet state are given by eq. 4, where the subscript i of \mathcal{Q}_1 shows the eigenvalue of s_Z .

 $\begin{aligned} & \mathcal{Q}_2 = \alpha\alpha\alpha\alpha \\ & \mathcal{Q}_1 = (1/2) \left[\alpha\alpha \left(\alpha\beta + \beta\alpha \right) + \left(\alpha\beta + \beta\alpha \right) \alpha\alpha \right] \\ & \mathcal{Q}_o = (1/\sqrt{6}) \left[\alpha\alpha\beta\beta + \left(\alpha\beta + \beta\alpha \right) \left(\alpha\beta + \beta\alpha \right) + \beta\beta\alpha\alpha \right] \\ & \mathcal{Q}_{-1} = (1/2) \left[\left(\alpha\beta + \beta\alpha \right)\beta\beta + \beta\beta \left(\alpha\beta + \beta\alpha \right) \right] \\ & \mathcal{Q}_{-2} = \beta\beta\beta\beta \end{aligned} \tag{4}$

The state of eigenvalue 2J is three-fold degenerate (eq. 5) with a total spin S=1.

$$\begin{split} &T_1 = (1/2) \left[\alpha \alpha (\alpha \beta + \beta \alpha) - (\alpha \beta + \beta \alpha) \alpha \alpha \right] \\ &T_0 = (1/\sqrt{2}) \left(-\alpha \alpha \beta \beta + \beta \beta \alpha \alpha \right) \\ &T_{-1} = (1/2) \left[(\alpha \beta + \beta \alpha) \beta \beta - \beta \beta (\alpha \beta + \beta \alpha) \right] \end{split} \tag{5}$$

The eigenfunction which belongs to the eigenvalue 4J is a singlet state (S=0) and given by eq. 6.

$$S_0 = (1/\sqrt{3}) \left[\alpha \alpha \beta \beta - (1/2) (\alpha \beta + \beta \alpha) (\alpha \beta + \beta \alpha) + \beta \beta \alpha \alpha \right]$$
 (6)

Since the valence-bond exchange integral J^{AB} is usually negative between open-shell organic molecules at a distance of the van der Waals contact, the ground state tends to become singlet as in Fig. 1b. In order to develop macroscopic ferromagnetism, it is highly necessary to find the conditions under which case Fig. 1a could result; the intermolecular interaction between high-spin molecules has to be ferromagnetic. How can this be designed? There is one such prediction by McConnell (ref. 8). According to this theory, an effective exchange interaction between two free radicals can be ferromagnetic when the product of spin densities $\rho_{\bf i}$ and $\rho_{\bf j}$ at two interacting sites i and j on different molecules A and B is negative in sign, since the valence-bond exchange integral $J_{\bf ij}$ is negative.(eq. 7).

$$\hat{\mathbf{H}}^{AB} = -2 \sum_{i,j} J_{i,j}^{AB} \mathbf{s}_{i}^{A} \cdot \mathbf{s}_{j}^{B} = -2 \mathbf{s}^{A} \cdot \mathbf{s}_{i,j}^{B} J_{i,j}^{AB} \rho_{i}^{A} \rho_{j}^{B}$$

$$(7)$$

Whereas the theory has been referred to interpreting complicated magnetic interactions found in crystals of stable organic radicals (ref. 9), it has never been tested using a logical model system.

Let us take parallel stacking of the two benzene rings one from each triplet diphenylcarbene molecule. Since the π -spin distribution in the latter is as shown in Fig. 2 (ref. 10), the

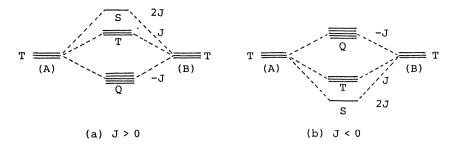


Fig. 1. Interaction diagrams for two triplet species A and B (ref. 11)

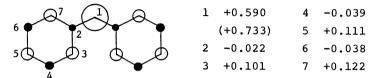


Fig. 2 π -Spin distribution in diphenylcarbene (ref. 10)

sign of the $\rho_1\rho_1$ at each interacting site between the two benzene rings is all negative in the ortho and para stacking patterns and all positive in the gem and meta (Fig. 3). Therefore, the ortho and para overlaps are predicted to give quintet ground states, while the gem and meta would lead to the ground singlet state. As a model system to test this theory, we have taken advantage of the [2.2]paracyclophane skeleton. Among the three isomers of bis-(phenylmethylenyl)[2.2]paracyclophanes $3\sim5$, the pseudoortho and pseudopara isomers satisfy the McConnell's condition to give the quintet ground state.

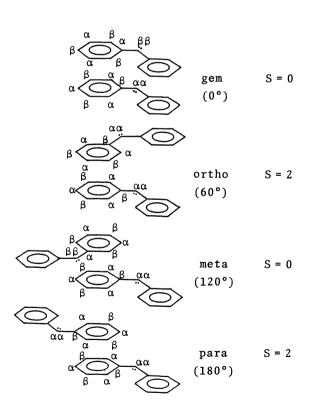
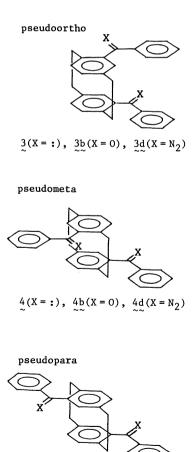


Fig. 3. Stacking patterns of the two benzene rings one from each diphenylcarbene molecule



 $5(X = :), 5b(X = 0), 5d(X = N_2)$

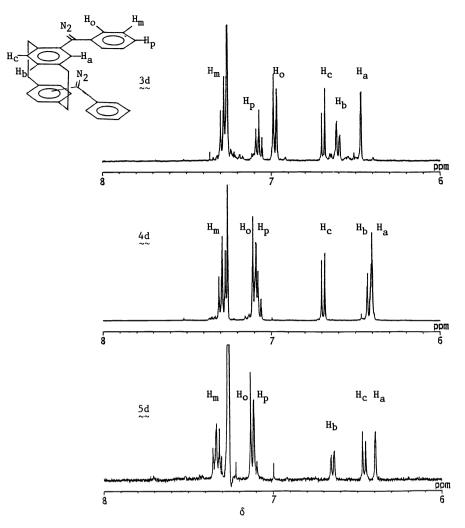


Fig. 4. Partial 400 MHz ¹H NMR spectra of didiazo compounds 3d, 4d and 5d

PREPARATION OF ISOMERIC BIS (α-DIAZOBENZYL)[2.2]PARACYCLOPHANES

Isomeric dibromo[2,2]paracyclophanes $3a \sim 5a$ were obtained by bromination of [2.2]paracyclophane (ref. 12). The dilithio compounds obtained from $3a \sim 5a$ on treatment with n-butyl-lithium in ether were allowed to react with benzaldehyde to give the corresponding diols. The psudoortho and pseudometa diols were oxidized with pyridinium chlorochromate in CH2Cl2 to give diketones 3b and 4b. Pseudopara diol was smoothly oxidized with sodium dichromate in acetic acid. Ketones $3b \sim 5b$ were resistant to hydrazine hydrate under typical conditions. The corresponding dihydrazones $3c \sim 5c$ were obtained by the reaction with excess anhydrous hydrazine and hydrazine monohydrochloride in DMSO at 110 °C for 10 h. Didiazo compounds $3d \sim 5d$ were obtained by oxidation of $3c \sim 5c$ with yellow HgO and purified by short-column chromatography on alumina: pseudoortho, mp 107-109 °C dec; pseudometa, mp 136-138 °C dec; pseudopara, mp 133-135 °C dec. Their 1 H NMR spectra are shown in Fig. 4.

ESR DETECTION OF ISOMERIC BIS(PHENYLMETHYLENYL)[2.2] PARACYCLOPHANES AND DETERMINATION OF THEIR SPIN STATES

Pseudoortho-bis(phenylmethylenyl)[2.2]paracyclophane

A solid solution of 3d in 2-methyltetrahydrofuran (2-MTHF) was irradiated with Pyrex-filtered light in an ESR cavity at 11 K. A set of intense ESR signals were obtained as shown in Fig. 5a. The spectrum as a whole resembles that of m-phenylenebis(phenylmethylene)($\frac{1}{2}$ (n = 2)) in the quintet ground state. The resonance positions and intensities of the signals are

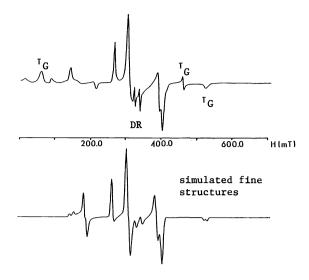


Fig. 5. ESR spectrum obtained on a Varian E-112/V7800 spectrometer with an X-band microwave unit by irradiating pseudoortho didiazo compound 3d in 2-MTHF at 11 K. Signals T and R are assigned to a monocarbene and \tilde{a} doublet radical, respectively, formed adventitiously. See text for the simulated spectrum.

reproduced reasonably well by a second-order perturbational calculation as a quintet species with zero-field splitting parameters $|D|=0.0624~\rm cm^{-1}$ and $|E|=0.0190~\rm cm^{-1}$ (Fig. 5b). The fitting between the observed and calculated spectra is not satisfactory for transitions at the lower magnetic field than 200 mT. This is due to the limitation of the high-field approximation used in the calculation (ref. 13). In addition to these quintet signals, triplet signals ($T_{\rm G}$) were detected at 62.8, 463.0, 526.7 and 732.0 mT and assigned to a monocarbene derived from $\frac{3d}{3d}$ ($|D|=0.3730~\rm cm^{-1}$ and $|E|=0.0156~\rm cm^{-1}$). Another triplet species with small zero-field splitting parameters ($|D|=0.0134~\rm cm^{-1}$ and $|E|=0.000~\rm cm^{-1}$) was detected in the central region of the spectrum. This is assigned to a triplet diradical (DR) formed presumably from dicarbene 3.

In order to determine the ground state spin multiplicity and thermal stability of the observed species, the temperature dependence of the signal intensities was studied (Fig. 6). The signal intensities of the quintet (Q) were found to obey Curie law in the temperature range $11 \sim 50$ K, indicating that the pseudoortho dicarbene 3 has the quintet ground state. This is the first experimental demonstration that the magnetic interaction between two diphenylcarbene moieties becomes ferromagnetic when two spin-containing benzene rings one from each diphenylcarbene are stacked in a proper mode dictated by McConnell (ref. 8). The signals due to $T_{\rm G}$ and DR also obeyed the Curie law in the temperature ranges $11 \sim 85$ K and $11 \sim 50$ K, respectively. At above 50 K, the latter signal intensities increased irreversibly at the expense of Q species, suggesting generation of DR from Q.

At temperatures higher than 20 K, a new signal started to appear at 104.0 mT. The intensity of the signal increased as the temperature was elevated and reached a maximum at 55 K (Fig.6). The signal decayed irreversibly at temperatures higher than 50 K. The thermal behavior is characteristic of a thermally populated species (T_T). Since relation $D_T = -3D_Q$ holds for two equivalent triplet species interacting weakly (ref. 11), the zero-field splitting parameters of the thermally populated triplet are estimated to be $|D| = 0.1872 \text{ cm}^{-1}$ and $|E| = 0.0570 \text{ cm}^{-1}$. The transition field of H_{min} for T_T is calculated to be 104.9 mT from the estimated D and E values in good agreement with the observed value of 104.0 mT. Since the signals due to the other transitions of T_T were weak and overlapped with intense Q signals, they could not be identified precisely. In reference to Fig. 1a and by assuming a Boltzman distribution of electron spins among the three states, the temperature dependence of the signal intensity for T_T is given by eq. 8.

$$I(T_{\rm T}) = \frac{C \exp\left(-\Delta E_{\rm QT}/RT\right)}{5 + 3 \exp\left(-\Delta E_{\rm OT}/RT\right) + \exp\left(-3\Delta E_{\rm OT}/2RT\right)} \cdot \frac{1}{T} \tag{8}$$

where $\Delta E_{\rm QT}$ is the energy gap between the ground Q and the thermally populated T states. The data in Fig. 6 give $\Delta E_{\rm QT} = 61~{\rm cm}^{-1}$ (= 175 cal/mol) and J is calculated from eq. 2 to be +16 cm⁻¹.

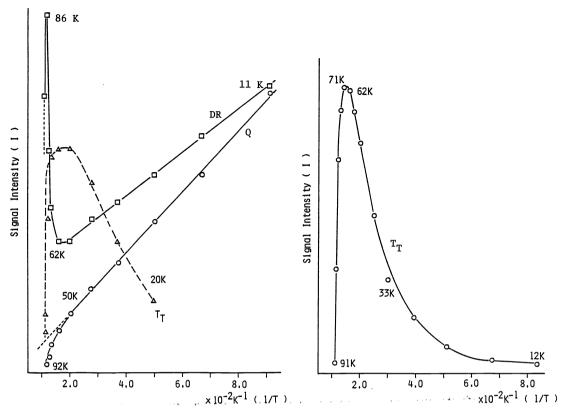


Fig. 6. Plots of ESR signal intensities vs the reciprocals of temperature for pseudoortho dicarbene 3.

Fig. 8. Plots of ESR signal intensities vs the reciprocals of temperature for pseudometa dicarbene $\frac{4}{2}$.

Pseudometa-bis(phenylmethylenyl)[2.2]paracyclophane

In sharp contrast with the pseudoortho isomer, signals due to the quintet species were not detected in the ESR spectrum obtained after irradiation of didiazo compound 4d at 11 K. Instead a set of signals due to a monocarbene ($|D|=0.3793~{\rm cm}^{-1}$ and $|E|=0.0158~{\rm cm}^{-1}$) were found. When the temperature of the irradiated solid solution was raised to 25 K, a new signal started to appear at 119.4 mT and its intensity increased with temperature giving a maximum at 65 K. Two additional peaks at 209.1 and 425.4 mT were discernible at this temperature (Fig. 7). The resonance fields of the three signals correspond to a triplet species (T_T) with $|D|=0.184~{\rm cm}^{-1}$ and $|E|=0.0023~{\rm cm}^{-1}$. Their thermal behavior (Fig. 8) is consistent with a thermally populated species.

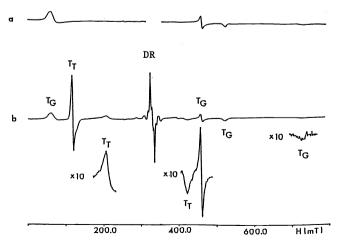


Fig. 7. ESR spectra obtained (a) in the dark after generation of pseudometa dicarbene 4 at 11 K, and (b) when sample (a) was warmed to 60 K. DR is due to a dirad \tilde{i} cal formed adventitiously

Since quintet signals were not detected and the observed triplet is thermally populated, the two diphenylcarbene units in 4 are concluded to interact antiferromagnetically to give a singlet ground state. The energy difference (ΔE_{ST}) between the ground singlet and higherlying triplet states (see Fig. 1b) is given by eq. 9:

$$I(T_{\rm T}) = \frac{\text{C} \exp{\left(-\Delta E_{\rm ST}/RT\right)}}{1 + 3\exp{\left(-\Delta E_{\rm ST}/RT\right)} + 5\exp{\left(-3\Delta E_{\rm ST}/RT\right)}} \cdot \frac{1}{T} \tag{9}$$

From the thermal behavior of the signal intensities of T_T in the temperature range 11 ~65 K, the $\Delta E_{\rm ST}$ value is estimated to be 98 cm⁻¹ (= 280 cal/mol), the J value being -49 cm⁻¹. The quintet state appears to be too high-lying to detect by thermal population.

Pseudopara-bis(phenylmethylenyl)[2.2]paracyclophane

Pseudopara didiazo compound 5d was photolyzed similarly in an ESR cavity at 15 K. Strong quintet signals due to dicarbene 5 were obtained: |D| = 0.1213 cm⁻¹ and |E| = 0.0063 cm⁻¹ (Fig. 9). In addition to these signals, multiplet signals due to a diradical species were observed at ca. 330 mT with the zero-field splitting parameters of |D| = 0.0110 cm⁻¹ and |E| = 0.00 cm⁻¹. Additional signals at 464.4, 530.1 and 742.4 mT are again attributed to a triplet monocarbene $(T_G)(|D| = 0.3837$ cm⁻¹ and |E| = 0.0168 cm⁻¹). The intensity of the quintet signals decreased as the temperature was raised from 15 to 41 K obeying apparently Curie law. However, the intensity was not recovered when the temperature was lowered from 41 K back to 15 K; about 60% of the original intensity was lost presumably by chemical reactions. The intensity of the diradical species increased during the annealing, showing that the quintet species may not be kinetically stable and converted to the diradical species even at these temperatures. Although it was not possible to obtain thermodynamical information from the temperature dependence of the signal intensities, the observation of the strong quintet signals strongly suggests that the quintet must be the ground state of dicarbene 5. A thermally populated triplet species was not detected in the temperature range $15 \sim 41$ K. The energy gap ΔE_{OT} in 5 appears to be even larger than that in 3.

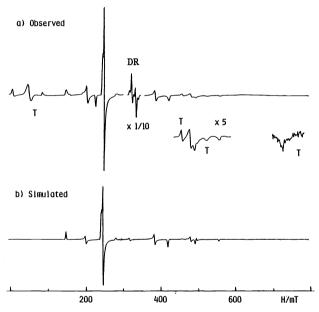


Fig. 9. ESR spectrum obtained by irradiation of pseudopara didiazo compound 5d in 2-MTHF at 11 K. $^{\sim}$

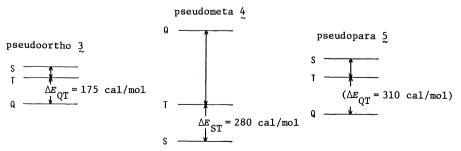


Fig. 10. The spin states of isomeric dicarbenes $3 \sim 5$.

The difference in the spin states of isomeric [2.2] paracyclophane dicarbenes $3 \, \bar{} 5$ is summarized in Fig. 10. The results serve as the first operational test for the McConnell's theory and lead to the conclusion that the mode of stacking of spin-containing benzene rings can be crucial in increasing the dimension of the intermolecular magnetic interaction of high-spin aromatic molecules in molecular assemblies.

MAGNETIC INTERACTION AMONG DIPHENYLCARBENE MOLECULES GENERATED IN CRYSTALS OF CORRESPONDING DIAZODIPHENYLMETHANES

The above findings have encouraged us to investigate intermolecular magnetic interaction among diphenylcarbene derivatives generated in the host crystals of the diazo precursors. If favorable molecular interactions as the ortho and para stacking in Fig. 3 are realized in host crystals, intermolecular spin alignments will be established. We have scrutinized a series of substituted diazodiphenylmethanes and found two extreme examples (ref. 14).

p-Methoxydiphenyldiazomethane, bis(p-methoxyphenyl)diazomethane and bis (p-octyloxyphenyl)diazomethane

When microcrystalline samples of p-methoxydiphenyldiazomethane 6a and bis(p-methoxyphenyl)-diazomethane 7a were irradiated in an ESR cavity at 10 K, the observed spectra were entirely different from those of the triplet obtained in glassy matrices (Fig. 11), but resembled more closely that of the quintet dicarbene 5 (Fig. 9). Intense signals at 248.9 and 259.9 mT for 6a and 247.6 and 259.0 mT for 7a were observed immediately after the start of irradiation. These lines are characteristic of quintet species (Q): $|D| = 0.136 \pm 0.005$ and $|E| = 0.013 \pm 0.004$ cm⁻¹ for 6a, and $|D| = 0.136 \pm 0.005$ and $|E| = 0.015 \pm 0.003$ cm⁻¹ for 7a. The observation of the intense Q signals suggests that pairs of triplet carbenes 6 and 7a generated in the crystals of 6a and 7a, respectively, interact efficiently to give the Q species. Weak triplet lines were also observed in the photolysates of 6a and 7a with the zero-field parameters of |D| = 0.408, |E| = 0.0180 cm⁻¹, and |D| = 0.408, |E| = 0.0194 cm⁻¹, respectively. These data are in good agreement with those obtained indisputably in 2-methyltetrahydrofuran glasses. Besides these lines, multiplet signals (M) were detected in the range $300 \sim 380$ mT (Fig. 11). The species responsible for M is estimated to be molecular aggregates between the Q pairs. Thus the ESR spectral feature appears to be well-explained by postulating that various sizes of clusters of the triplet carbene species are formed by photolysis of the diazo compounds in crystals.

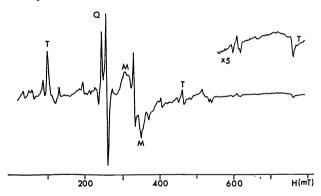


Fig. 11. ESR spectrum obtained by irradiation of the microcrystalline sample of bis(p-methoxyphenyl)diazomethane at 10 K.

Temperature dependence of the intensities of the M, Q and T signals was studied in the range $14 \, ^{\sim} \, 120$ K. The Q species was kinetically stable up to 60 K and disappeared rapidly at 80 K. The Curie plot of the Q signals shows a concave departure from linearity, and is partly reproduced by a model in which Q is the ground state and a thermally populated triplet lies about 30 cal/mol above the ground state. Most of the T signals observed is then reasoned to be due to the triplet carbene 7 isolated magnetically in the crystal.

Since the energy of the interaction between electron spins and an external magnetic field is smaller than the thermal energy, orientation of the spins, i.e., strength of magnetization, increases with that of the external field and decreases with temperature as dictated by the Brillouin function $B_J(x)$:

$$M = NgJ\mu_{\mathbf{R}}B_{\mathbf{J}}(\mathbf{x}) \tag{10}$$

where

$$B_{\mathbf{J}}(\mathbf{x}) = \frac{2J+1}{2J} \; \coth(\frac{2J+1}{2J} \; \mathbf{x}) \; -\frac{1}{2J} \; \coth(\frac{\mathbf{x}}{2J}) \qquad \text{and} \; \; X = gJ\mu_{\mathbf{B}}H/(k_{\mathbf{B}}T)$$

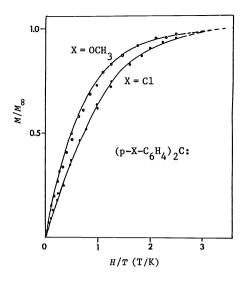


Fig. 12. Plots of magnetization of bis(p-methoxyphenyl)carbene (o) and bis(p-chlorophenyl)carbene (•) generated in the corresponding diazo crystals vs temperaturenormalized field strength(ref. 14).

At very strong external magnetic field and very low temperature $(H/T \sim \infty)$, all the spins will be oriented to the direction of the magnetic field and saturation of magnetization will occur. When normalized to this value M_{∞} , the magnetization curve should be equal to $B_{\rm J}(x)$. Such plots for bis(p-methoxyphenyl)carbene ${\it I}$ are shown in Fig. 12. Since ${\it I}$ is composed of light atoms with small spin-orbit coupling, ${\it J}=S$ and ${\it g}=2$. The experimental data are found by a non-linear least-square method to fit the theoretical curve with ${\it S}=(3.8\pm0.4)/2$, providing independent evidence for the formation of a pair of triplet 7 coupled ferromagnetically.

The ferromagnetic coupling of \S and \S generated in the precursor crystals of \S a and \S a, respectively, is expected to arise from the intermolecular exchange interaction that is effectively positive. The structure of the host crystals should be informative. Since it was difficult to obtain good single crystals amenable to an X-ray analysis, we studied the crystal structure $(P2_1/a)$ of isomorphous p,p'-dimethoxybenzophenone (ref. 16). Among four independent molecules in the unit cell, one benzene ring of the ketone molecule is nearly in parallel with the benzene ring of an adjacent molecule, taking a pseudopara-type orientation (see Fig. 3). Strictly speaking, the centers of the two interacting benzene rings are slightly off and only two carbon atoms on one benzene ring are interacting effectively, but the stacking pattern satisfies McConnell's theory in a sense that the carbon atoms with positive spin density face those with negative ones at the interacting sites. The molecular packing in the crystal also rationalizes the presence of weak interaction between two carbone pairs each in the quintet state.

In order to design an assembly of carbene molecules that have intermolecular spin alignment over several molecules, we have taken advantage of the dispersion force of long alkyl chains (ref. 17). The ESR spectrum of the photolysate of polycrystalline bis(p-octyloxyphenyl)-diazomethane 8a at 10 K showed a complex multiplet in the range 260° 380 mT, accompanied by weak quintet signals at 245.8 and 258 mT. A plot of the inverse of the paramagnetic susceptibility vs temperature in the range 2° 150 K consisted of four linear regions, showing development of various molecular fields at different temperatures (ref. 15). From the magnetization behavior at 2.1 K, S=8/2 was obtained. Intermolecular magnetic interaction now develops over four carbene molecules.

Bis(p-chlorophenyl)diazomethane

When irradiated at 10 K, monoclinic crystals of bis(p-chlorophenyl)diazomethane showed strong triplet signals with the zero-field parameters (|D| = 0.3980 and |E| = 0.0194 cm⁻¹) very close to those obtained in 2-MTHF glasses. A very weak signal assigned to a quintet was detected at 280 mT. From an analysis of the temperature dependence of these signals similar to those carried out for the methoxy derivatives, the carbene pair was found to be in the ground singlet state with the triplet and quintet states thermally populated. An X-ray analysis has revealed that the molecular packing is like the geminal overlap in Fig. 3 (ref. 14).

CONCLUSION

We have been able to demonstrate by model experiments that there are unique stacking patterns of the spin-distributed benzene rings of diphenylcarbene units leading to intermolecular electron spin alignment. They are ortho and para stacking as dictated by the McConnell's theory. When oriented by the dispersion force of long alkyl chains in the crystal of a diazo precursor, the spin alignment over four triplet molecules was achieved.

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Previously we showed that a large number of electron spins can be aligned systematically in organic molecules when properly designed as in polycarbenes 1 and 2. We are therefore in the position to take advantage of these guiding principles to develop highly spin-aligned organic molecular systems that should have interesting magnetic properties.

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